

The main ensembles: Canonical

①

Let's first assume we have a fixed number of particles $N \rightarrow$ absolute knowledge about this part of the problem

$$\mathcal{D}_B = \frac{1}{Z} \exp -\beta \hat{H}$$

Let's now use the basis that diagonalizes \hat{H}

$$\hat{H} |r\rangle = E_r |r\rangle$$

$$\Rightarrow \langle r | \mathcal{D}_B | r' \rangle = \frac{1}{Z} \langle r | \underbrace{\exp -\beta \hat{H}}_{f(H)} | r' \rangle = \frac{1}{Z} \exp -\beta E_r \delta_{rr'}$$

$$\mathcal{D}_B = \frac{1}{Z} \sum_r e^{-\beta E_r} |r\rangle \langle r|$$

if $H = H^{(a)} + H^{(b)}$ with $[H^{(a)}, H^{(b)}] = 0$

$$\Rightarrow \mathcal{D}_B = \mathcal{D}_B^{(a)} \otimes \mathcal{D}_B^{(b)}$$

Let $|l\rangle = |a\rangle \otimes |\alpha\rangle$ ← eigenstate of $H^{(a)}$ ②
↑
eigenstate of $H^{(a)}$
 $|m\rangle = |b\rangle \otimes |\beta\rangle$ ←

$$\langle l | \hat{H} | m \rangle = (H_{ab}^{(a)} + H_{\alpha\beta}^{(a)}) \delta_{ab} \delta_{\alpha\beta}$$

The same way

$$\langle l | e^{-\beta \hat{H}} | m \rangle = \left(e^{-\beta H_{ab}^{(a)}} e^{-\beta H_{\alpha\beta}^{(a)}} \right) \delta_{ab} \delta_{\alpha\beta}$$

True only if they commute

Campbell-Baker-Hausdorff formula

Satisfied if the Hamiltonian is separable

$$\begin{aligned} Z &= \text{Tr} e^{-\beta \hat{H}} = \sum_l \langle l | e^{-\beta \hat{H}} | l \rangle = \sum_{a\alpha} \langle a | \otimes \langle \alpha | e^{-\beta \hat{H}} | a \rangle \otimes | \alpha \rangle \\ &= \sum_{a\alpha} e^{-\beta H_{aa}^{(a)}} e^{-\beta H_{\alpha\alpha}^{(a)}} = \sum_a e^{-\beta H_{aa}^{(a)}} \sum_{\alpha} e^{-\beta H_{\alpha\alpha}^{(a)}} \end{aligned}$$

$$\Rightarrow Z = Z^{(a)} Z^{(\alpha)}$$

Canonical ensemble

$$Z = \text{Tr} e^{-\beta \hat{H}} = \sum_r e^{-\beta E_r}$$

In the continuum limit $\sum_r \rightarrow \int dE \rho(E)$

$$Z = \int dE \rho(E) e^{-\beta E}$$

$$\left[E(\beta, \alpha) = - \frac{\partial \ln Z}{\partial \alpha_i} \right]_{\beta, \alpha_j \neq \alpha_i} \quad (3)$$

Let's also call

$$\begin{aligned} X_i &= \left\langle \frac{\partial H}{\partial \alpha_i} \right\rangle = \text{Tr} \left[\frac{\partial H}{\partial \alpha_i} D \right] = \text{Tr} \left[\frac{1}{Z} \frac{\partial H}{\partial \alpha_i} e^{-\beta H} \right] \\ &= \text{Tr} \left[- \frac{1}{Z} \frac{1}{\beta} \frac{\partial}{\partial \alpha_i} e^{-\beta H} \right] = - \frac{1}{\beta} \frac{1}{Z} \text{Tr} \left[\frac{\partial}{\partial \alpha_i} e^{-\beta H} \right] \\ &= - \frac{1}{\beta} \frac{1}{Z} \frac{\partial}{\partial \alpha_i} \underbrace{\left[\text{Tr} e^{-\beta H} \right]}_Z = - \frac{1}{\beta} \frac{1}{Z} \frac{\partial Z}{\partial \alpha_i} \end{aligned}$$

$$\Rightarrow X_i = \left\langle \frac{\partial H}{\partial \alpha_i} \right\rangle = - \frac{1}{\beta} \frac{\partial \ln Z}{\partial \alpha_i} \Big|_{\beta, \alpha_j \neq \alpha_i}$$

Note: The \hat{A}_i associated with Lagrange multiplier α_i have the form $\langle A_i \rangle \propto \frac{\partial \ln Z}{\partial \alpha_i}$

But also other observables will also be of the form $\langle \hat{O} \rangle \propto \frac{\partial \ln Z}{\partial x}$

where x is some parameter

In general

$$\rightarrow \langle H \rangle = E$$

(4)

$$S = k \ln Z - k \sum_i \lambda_i A_i$$

$$S = k \ln Z + \beta k E$$

Reorganizing

$$k \ln Z = S - \beta k E = S - \left(\frac{1}{T}\right) E$$

We note that

$$S(E, V)$$

$$\text{and } \frac{\partial S}{\partial E} = \frac{1}{T}$$

Legendre transform

$$f(x) \quad \frac{\partial f}{\partial x} = u$$

$$\Rightarrow g(u) = f(x) - ux$$

$$\Rightarrow \begin{aligned} f(x) &\rightarrow S(E, V) \\ u &\rightarrow \frac{\partial f(x)}{\partial x} \rightarrow \frac{1}{T} \\ x &\rightarrow E \\ g(u) &\rightarrow \Phi_1\left(\frac{1}{T}, V\right) \end{aligned}$$

$$\ln Z(\beta, V) = \frac{1}{k} \Phi_1\left(\frac{1}{T}, V\right)$$

Rearranging again

$$-kT \ln Z = \underbrace{E - TS}_{\text{Legendre transform}} = F(T, V)$$

Legendre transform

$$\Rightarrow F(T, V) = -\frac{1}{\beta} \ln Z$$

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In differential form

$$dF = dE - T dS - S dT + P dV$$

$$\frac{1}{k} dS = d(\ln Z) + \dots$$

$$Z = Z[\beta, x_i]$$

$$\Rightarrow d \ln Z = \underbrace{\frac{\partial \ln Z}{\partial \beta}}_{-E} d\beta + \sum_i \underbrace{\frac{\partial \ln Z}{\partial x_i}}_{-\beta X_i} dx_i$$

$$= -E d\beta - \beta \sum_i X_i dx_i$$

$$d(\beta E) = E d\beta + \beta dE$$

$$\Rightarrow \frac{1}{k} dS = \cancel{-E d\beta} - \beta \sum_i X_i dx_i + \cancel{E d\beta} + \beta dE$$

$$\frac{1}{k} dS = \beta dE - \beta \sum_i X_i dx_i$$

$$\textcircled{1} \quad T dS = dE - \sum_i X_i dx_i$$

Also

$$dF = -k_B \ln Z dT - k_B T d \ln Z$$

$$\frac{S}{k} - \beta E \quad -E d\beta - \beta \sum_i X_i dx_i$$

$$= -S dT + \cancel{\frac{E}{T} dT} + \frac{E}{\beta} d\beta + \sum_i X_i dx_i \quad \textcircled{2}$$

$$\beta = \frac{1}{kT} \quad d\beta = -\frac{1}{kT^2} dT$$

$$\downarrow d\beta = \frac{\beta}{T}$$

$$\textcircled{2} \quad dF = -SdT + \sum_i X_i dx_i$$

From our thermodynamic relation

$$dE = TdS - PdV$$

$$\textcircled{1} \quad dE = TdS - X_i dV$$

$$\Rightarrow x_i = -V \quad X_i = P$$

$$X_i = \left\langle \frac{\partial H}{\partial x_i} \right\rangle = - \left\langle \frac{\partial H}{\partial V} \right\rangle \equiv \text{units of pressure}$$

consistent

From thermodynamics

$$\textcircled{2} \quad \frac{\partial F}{\partial V} = -P$$

$$\left. \frac{\partial F}{\partial V} \right|_{T, N} = -P$$

consistent

$$\Rightarrow P = \frac{1}{\beta} \frac{\partial \ln Z}{\partial V} \Big|_{\beta, N}$$

$$\left. \frac{\partial F}{\partial N} \right|_{T, V} = \mu$$

$$\mu = \frac{1}{\beta} \frac{\partial \ln Z}{\partial N} \Big|_{\beta, V}$$

Fluctuations

$$-\frac{\partial E}{\partial \beta} = \int d^2 \text{Tr} \left[(\hat{H} - E) \mathcal{D}^2 (\hat{H} - E) \mathcal{D}^{(1-2)} \right] = \text{Tr} \left[(\hat{H} - E) \mathcal{D} \right]$$

$\partial \beta$ \circ L

$$D = \frac{1}{Z} e^{-\beta H} \quad \downarrow \checkmark$$

\Rightarrow commutes with H

$$= \langle (\hat{H} - E)^2 \rangle = \langle \hat{H}^2 \rangle - E^2$$

$$E = - \frac{\partial \ln Z}{\partial \beta}$$

$$\langle (\hat{H} - E)^2 \rangle = \frac{\partial^2 \ln Z}{\partial \beta^2}$$

The heat capacity

$$C_x = T \frac{\partial S}{\partial T} \Big|_x = \text{heat capacity at constant } x$$

$$\text{but } \frac{\partial S}{\partial T} \Big|_x = \underbrace{\frac{\partial S}{\partial E}}_{\frac{1}{T}} \Big|_x \frac{\partial E}{\partial T} \Big|_x \Rightarrow C_x = \frac{\partial E}{\partial T} \Big|_x$$

$$\frac{\partial E}{\partial T} \Big|_x = \frac{\partial E}{\partial \beta} \Big|_x \frac{\partial \beta}{\partial T} \Big|_x = - \frac{1}{KT^2} \frac{\partial E}{\partial \beta}$$

but

$$\frac{\partial E}{\partial \beta} = - \langle (\hat{H} - E)^2 \rangle$$

$$C_x = \frac{\langle (\hat{H} - E)^2 \rangle}{KT^2} \geq 0$$

consistent

$\ln Z$ is convex

