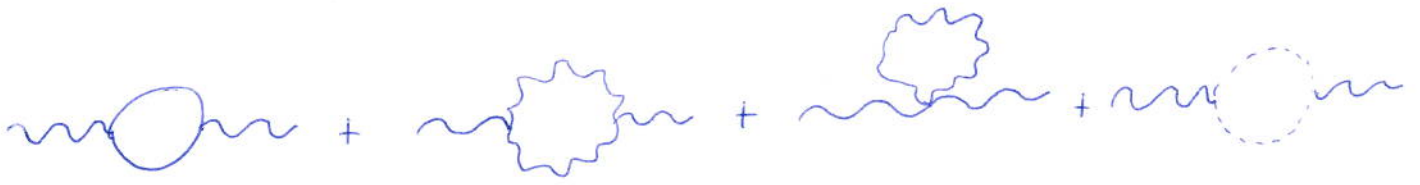


Gauge boson self-energy

$$q^\mu \left(\underbrace{\text{---}\bigcirc\text{---}} \right) = 0 \quad ; \text{ Ward-identity}$$

$$i(q^2 g^{\mu\nu} - q^\mu q^\nu) \Pi(q^2)$$

Contribution until g^2



$$= i(q^2 g^{\mu\nu} - q^\mu q^\nu) \delta^{ab} \left(\frac{-g^2}{(4\pi)^2} \frac{4}{3} n_f C(\alpha_2) \Gamma(2 - \frac{d}{2}) + \dots \right)$$

$$+ i(g^{\mu\nu} q^2 - q^\mu q^\nu) \delta^{ab} \left(\frac{-g^2}{(4\pi)^2} \underbrace{\left(-\frac{13}{6} + \frac{5}{2} \right)}_{-\frac{5}{3}} C_2(G) \Gamma(2 - \frac{d}{2}) \right) + \dots$$

$$\text{wavy line} = -i(k^2 g^{\mu\nu} - k^\mu k^\nu) \delta^{ab} \delta_3$$

$$\text{crossed lines} = i \not{k} \delta_2$$

$$\text{vertex} = i g t^a \gamma^\mu \delta_1$$

Gauge boson propagator:

$$G^{(2)}(p) = \text{wavy line} + (\text{Lead order loop}) + \text{wavy line}$$

$$\tilde{D}(k)_{\mu\nu}^{ab} = \frac{-i}{k^2 + i\epsilon} \left[g_{\mu\nu} + (\xi - 1) \frac{k_\mu k_\nu}{k^2} \right] \delta_{ab}$$

$$\begin{matrix} a & & b \\ \mu & \text{wavy line} & \nu \end{matrix}$$

(Lead Loop): $\tilde{D}(k)_{\mu\alpha}^{ac} \left[i(g^{\alpha\beta} k^2 - k^\alpha k^\beta) \delta^{cd} \right] \delta^3 \tilde{D}(k)_{\beta\nu}^{db}$

wavy line: $\tilde{D}(k)_{\mu\alpha}^{ac} \left[-i(k^2 g^{\alpha\beta} - k^\alpha k^\beta) \delta^{cd} \right] \delta^3 \tilde{D}(k)_{\beta\nu}^{db}$

$$\bullet \frac{k_\mu k_\alpha}{k^2} (k^2 g^{\alpha\beta} - k^\alpha k^\beta) = \frac{k_\mu k_\alpha}{k^2} (k^2 k^{\alpha\beta} - k^\alpha k^\beta) = 0$$

$$\begin{aligned} \bullet \tilde{D}(k)_{\mu\alpha}^{ac} (k^2 g^{\alpha\beta} - k^\alpha k^\beta) \delta^{cd} \tilde{D}(k)_{\beta\nu}^{db} &= \delta^{ac} \delta^{cd} \delta^{db} \left(-i \frac{g_{\mu\alpha}}{k^2} \right) \\ &\quad \times (k^2 g^{\alpha\beta} - k^\alpha k^\beta) \left(-i \frac{g_{\beta\nu}}{k^2} \right) \\ &= - \frac{\delta^{ab}}{k^4} (k^2 g_{\mu\nu} - k_\mu k_\nu) \end{aligned}$$

$$\frac{dG^{(2)}}{dM} = \left[\frac{i}{k^4} (k^2 g_{\mu\nu} - k_\mu k_\nu) \delta^{ab} \right] \frac{d}{dM} \delta^3$$

From Callan-Symanzik

$$M \frac{d}{dM} G^{(2)} + \beta(g) \frac{d}{dg} G^{(2)} + 2 \gamma(g) G^{(2)} = 0$$

Comparing in order of g

$$A = \mathcal{O}(g^2)$$

$$\underbrace{\beta(g)}_{\mathcal{O}(g^n)} \underbrace{\frac{d}{dg} G^{(2)}}_{\mathcal{O}(g)} + 2 \underbrace{\gamma(g)}_{\mathcal{O}(g^n)} \underbrace{G^{(2)}}_{\mathcal{O}(1)}$$

First order, ignore β

$$2 \gamma_3(g) G^{(2)} = -M \frac{d}{dM} G^{(2)}$$

$$2 \gamma_3(g) \left\{ \frac{i}{k^2} \left[g_{\mu\nu} - \frac{k_\mu k_\nu}{k^2} \right] \delta^{ab} \right\} = M \left\{ \frac{i}{k^4} (k^2 g_{\mu\nu} - k_\mu k_\nu) \delta^{ab} \right\}$$

$$\boxed{\gamma_3(g) = \frac{1}{2} M \frac{d}{dM} \delta^3}$$

$$\times \frac{d}{dM} \delta^3$$

Same for fermion

$$\gamma_2(g) = \frac{1}{2} M \frac{d}{dM} \delta_2$$

Green three-point function:

$$\begin{aligned}
 &= \left[\text{tree} + \text{1PI} + \text{2-loop} \right] \left[\text{prop} + \text{1PI-prop} \right] \\
 &\quad \times \prod_1^2 \left[\text{prop} + \text{1PI-prop} + \text{2-loop-prop} \right]
 \end{aligned}$$

From Callan-Symanzik:

$$\left[M \frac{d}{dM} + \beta(g) \frac{d}{dg} + 2\gamma_2(g) + \gamma_3(g) \right] G^{(3)} = 0$$

$$\begin{aligned}
 M \frac{d}{dM} G^{(3)} &= M \frac{d}{dM} \left[\underbrace{\left(\overset{ig\delta_1}{\text{1PI}} \right)}_{ig} (\text{prop}) (\text{prop})^2 + 2 \left(\overset{ig}{\text{1PI}} \right) (\text{prop}) (\text{prop}) \left(\overset{-\delta_2}{\text{2-loop-prop}} \right) \right. \\
 &\quad \left. + \left(\overset{ig}{\text{1PI}} \right) \underbrace{(\text{prop})}_{-\delta_3} (\text{prop})^2 \right]
 \end{aligned}$$

$$\bullet M \frac{d}{dM} G^{(3)} = (\text{wavy}) (\rightarrow)^2 M \frac{d}{dM} (ig \delta_1 - i2g \delta_2 - ig \delta_3)$$

$$\bullet \frac{d}{dg} G^{(3)} = i (\text{wavy}) (\rightarrow)^2$$

$$G^{(3)} = \overset{\text{L.O.}}{\nearrow} ig (\text{wavy}) (\rightarrow)^2$$

Replacing back into C-5.

$$(\text{wavy}) (\rightarrow)^2 \left[M \frac{d}{dM} (ig \delta_1 - i2g \delta_2 - ig \delta_3) + i\mathcal{B}(g) + ig M \frac{d}{dM} \delta_2 + ig \frac{1}{2} M \frac{d}{dM} \delta_3 \right] = 0$$

$$\mathcal{B}(g) = g M \frac{d}{dM} \left(-\delta_1 + 2\delta_2 + \delta_3 - \delta_2 - \frac{1}{2} \delta_3 \right)$$

$$\mathcal{B}(g) = g M \frac{d}{dM} \left[-\delta_1 + \delta_2 + \frac{1}{2} \delta_3 \right]$$

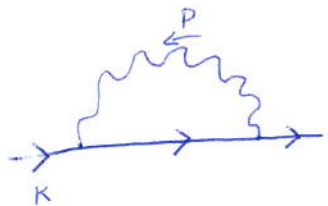
Determine counter terms:

$$\delta_3 = -\frac{g^2}{(4\pi)^2} \left[\frac{4}{3} n_f C(R) + \left(-\frac{5}{3}\right) C_2(G) \right]$$

should cancel the divergence from Gauge boson self energy (order g^2)

$$\delta_3 = \frac{g^2}{(4\pi)^2} \left[\frac{5}{3} C_2(G) - \frac{4}{3} n_f C(R) \right] \frac{\Gamma(2-d/2)}{(M^2)^{2-d/2}}$$

Now δ_2 :

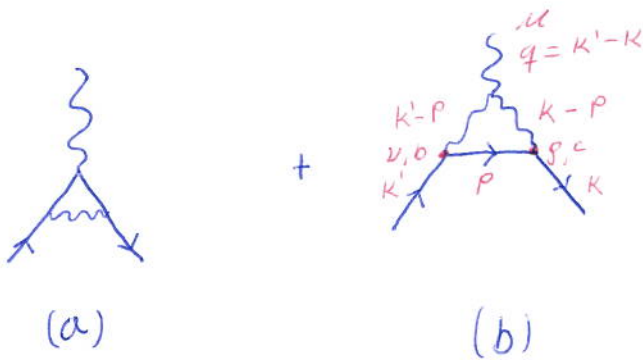


$$: \int \frac{d^4 p}{(2\pi)^4} (ig)^2 \gamma^\mu t^a \frac{i(\not{p} + \not{k})}{(p+k)^2} \gamma_\mu t^a \frac{-i}{p^2}$$

$$= \frac{ig^2}{(4\pi)^2} \underbrace{K C_2(R) \Gamma(2-d/2)}_{\text{divergent part}} + \dots$$

$$\delta_2 = -\frac{g^2}{(4\pi)^2} \frac{\Gamma(2-d/2)}{(M^2)^{2-d/2}} C_2(R)$$

Now d_1 , come from vertex correction 1-loop.



$$(a): \int \frac{d^4 p}{(2\pi)^4} g^3 t^b t^a t^b \frac{\gamma^\nu (\not{p} + \not{k}') \gamma^\nu (\not{p} + \not{k}) \gamma_\nu}{(p+k')^2 (p+k)^2 p^2}$$

using, gauge group matrices

$$\begin{aligned} t^b t^a t^b &= t^b t^b t^a + t^b [t^a, t^b] \\ &= C_2(R) t^a + i t^b f^{abc} t^c \\ &= C_2(R) t^a + i \frac{1}{2} [f^{abc} t^b t^c + f^{abc} t^a t^c] \\ &= C_2(R) t^a + \frac{1}{2} f^{abc} [t^b, t^c] \\ &= C_2(R) t^a + \frac{i}{2} f^{abc} f^{bcd} t^d \\ &= C_2(R) t^a - \frac{1}{2} f^{abc} f^{bcd} t^d \\ &= \left[C_2(R) - \frac{1}{2} C_2(G) \right] t^a \end{aligned}$$

then, the divergent part for large P

$$\sim g^3 \left[C_2(R) - \frac{1}{2} C_2(G) \right] t^a \int \frac{d^4 p}{(2\pi)^4} \frac{\gamma^\nu \not{p} \gamma^\mu \not{p} \gamma_\nu}{(p^2)^3}$$

$$= \frac{i g^3}{(4\pi)^2} \left[C_2(R) - \frac{1}{2} C_2(G) \right] t^a \gamma^\mu \left[\Gamma(2 - d/2) + \dots \right]$$

(b): $\int \frac{d^4 p}{(2\pi)^4} (i g \gamma_\nu t^b) \frac{i \not{p}}{p^2} (i g \gamma_\rho t^c) \frac{-i}{(k'-p)^2} \frac{-i}{(k-p)^2}$

$$\times g f^{abc} \left[g^{\mu\nu} (q+k'-p)^\rho + g^{\nu\rho} (-k'-k+p)^\mu + g^{\rho\mu} (k-p-q)^\nu \right]$$

Note: $f^{abc} t^b t^c = \frac{1}{2} f^{abcd} t^d = \frac{i}{2} C_2(G) t^a$

then

$\underbrace{P \gg k', k, q}_{\text{PS}}$ $\frac{g^3}{2} C_2(G) t^a \int \frac{d^4 p}{(2\pi)^4} \gamma_\nu \not{p} \gamma_\rho \frac{g^{\mu\nu} p^\rho - 2 g^{\nu\rho} p^\mu + g^{\rho\mu} p^\nu}{(p^2)^3}$

$$\sim \frac{i g^3}{(4\pi)^2} \frac{3}{2} C_2(G) t^a \gamma^\mu \left(\Gamma(2 - d/2) + \dots \right)$$

the counter term should cancel this divergent part

$$\delta_1 = \frac{-g^2}{(4\pi)^2} \frac{\Gamma(2 - d/2)}{(M^2)^{2 - d/2}} \left[C_2(R) + C_2(G) \right]$$

$$\beta(g) = g M \frac{d}{dM} \left[-\mathcal{J}_1 + \mathcal{J}_2 + \frac{1}{2} \mathcal{J}_3 \right]$$

Note: $\bullet M \frac{d}{dM} \frac{1}{(M^2)^{2-d/2}} = (d-4) M^{d-4}$

$\bullet M^{d-4} \Gamma(2 - \frac{d}{2}) \sim -\frac{2}{d-4} + \dots$

$$\beta(g) = -2g \left[\frac{g^2}{(4\pi)^2} (C_2(R) + C_2(G)) - \frac{g^2}{(4\pi)^2} C_2(R) \right.$$

$$\left. + \frac{1}{2} \frac{g^2}{(4\pi)^2} \left(\frac{5}{3} C_2(G) - \frac{4}{3} n_f C(R) \right) \right]$$

$$= -\frac{g^3}{(4\pi)^2} \left[\frac{11}{3} C_2(G) - \frac{4}{3} n_f C(R) \right]$$

Negative for $n_f \cdot C(R) < \frac{11}{4} C_2(G)$

$$\left. \begin{array}{l} \text{SU}(N): C_2(G) = N \\ \text{fundamental } C(N) = \frac{1}{2} \\ \text{Fermions} \end{array} \right\} n_f < \frac{11}{2} N$$